



Atomic Resolution Holography Using Thermal Neutrons

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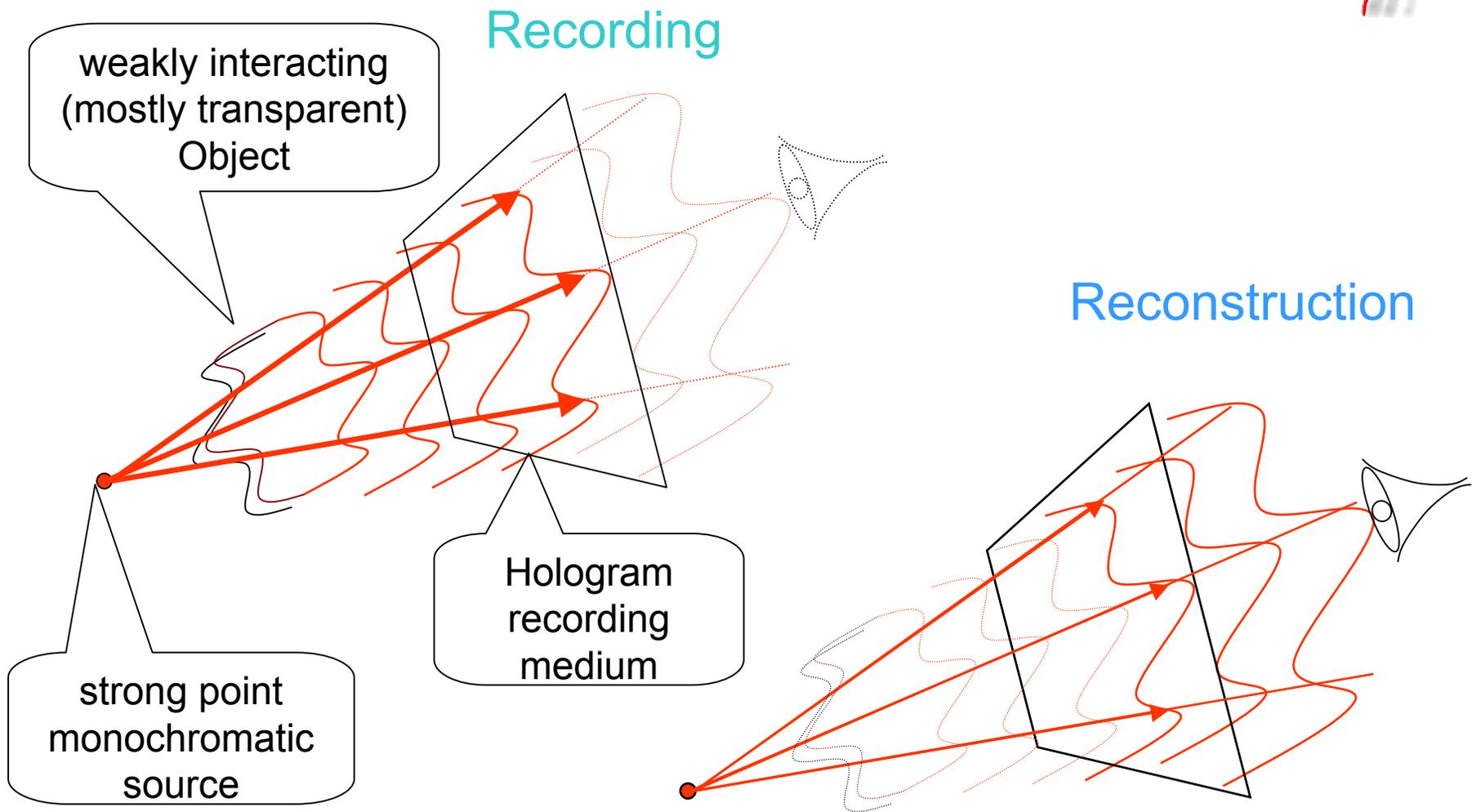
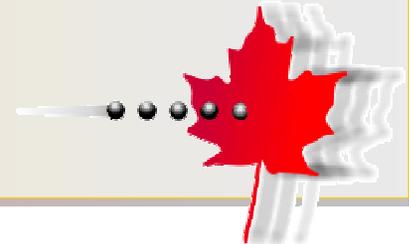
Bhaskar Sur

Vinicius Anghel

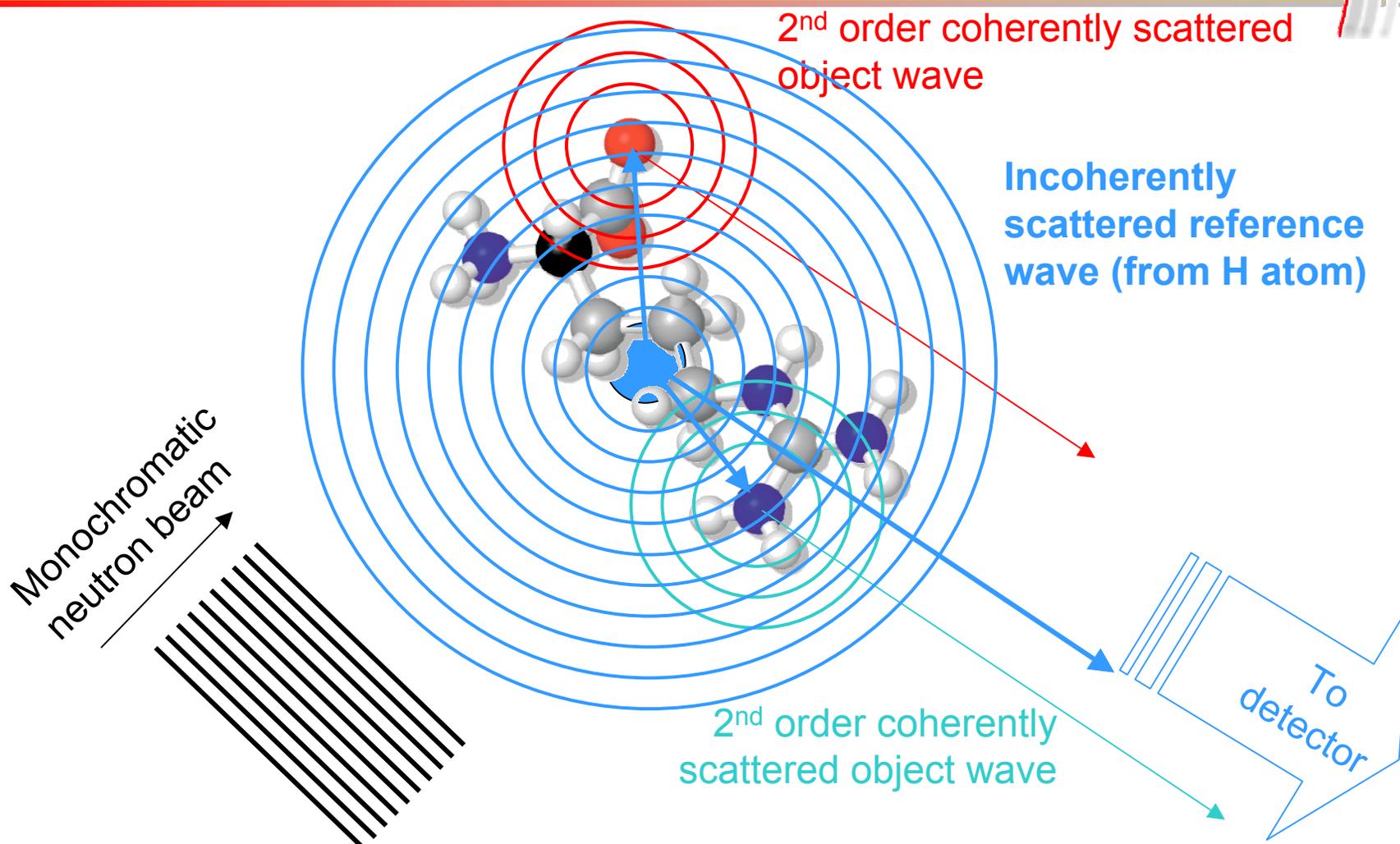
Atomic Energy of Canada Limited
Chalk River Laboratories



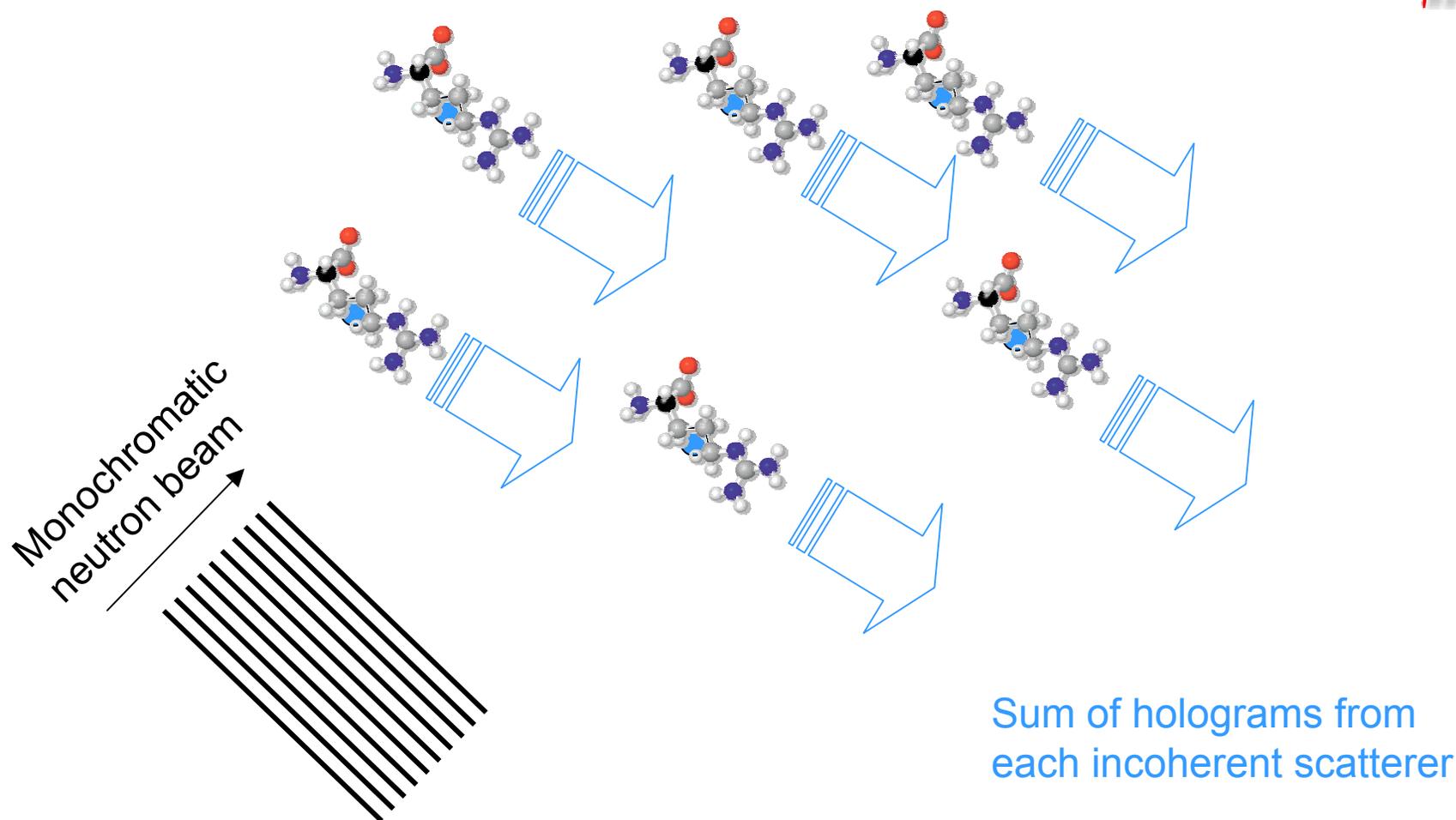
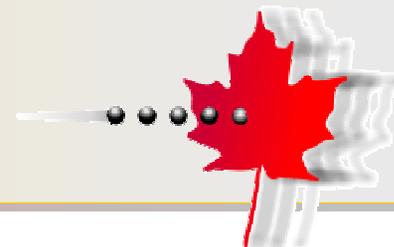
Gabor Holography



Atomic Structure Holography with thermal neutron waves



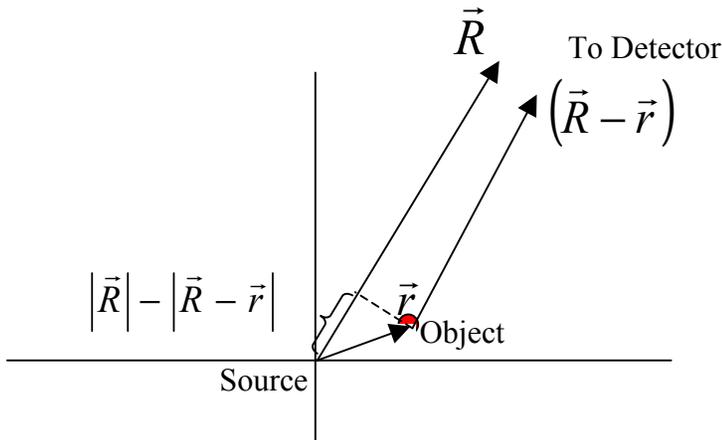
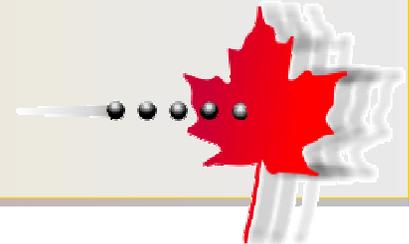
Atomic structure holography of macroscopic samples



Oriented (not necessarily crystalline) samples required

ARC - CMRC

S-Wave Diffraction for a Single Point Source



Convert detector coordinates to direction for $R \gg r$:

$$k(|\mathbf{R}| - |\mathbf{R} - \mathbf{r}_0|) = k(\mathbf{r}_0 \cdot \hat{\mathbf{R}}) = \mathbf{k} \cdot \mathbf{r}_0$$

$$\psi_{ref}(\mathbf{R}) = a \frac{e^{ikR}}{R}$$

Time-independent (stationary), scalar, spherical waves centered at $\mathbf{R} - \mathbf{r}$

$$\psi_{obj}(\mathbf{R}) = a \sum_{j=1}^{j=n} \frac{e^{ikr_j}}{r_j} b_j \frac{e^{ik|\mathbf{R}-\mathbf{r}_j|}}{|\mathbf{R}-\mathbf{r}_j|}$$

$$\psi_{tot}(\mathbf{R}) = a \frac{e^{ikR}}{R} \left(1 + \sum_{j=1}^{j=n} \left\{ \frac{R}{|\mathbf{R}-\mathbf{r}_j|} \right\} \frac{b_j}{r_j} e^{i(kr_j - k\{R - |\mathbf{R}-\mathbf{r}_j|\})} \right)$$

$$\chi(\mathbf{k}) = \sum_{j=1}^{j=n} \frac{b_j}{r_j} e^{i(kr_j - \mathbf{k} \cdot \mathbf{r}_j)} \quad \text{complex modulation function}$$

detected intensity $I(\mathbf{k}) = \psi_{tot}^*(\mathbf{R}) \psi_{tot}(\mathbf{R}) \approx \frac{a^* a}{R^2} \left(1 + 2\text{Real}\{\chi(\mathbf{k})\} + |\chi(\mathbf{k})|^2 \right)$

constant small much smaller

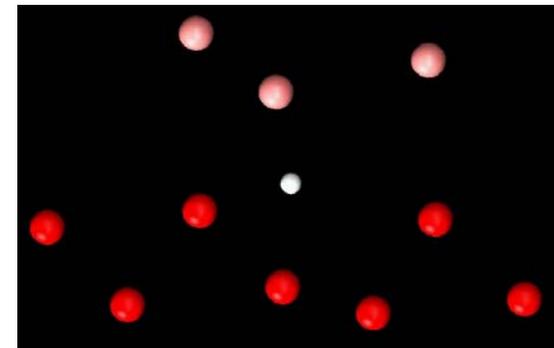
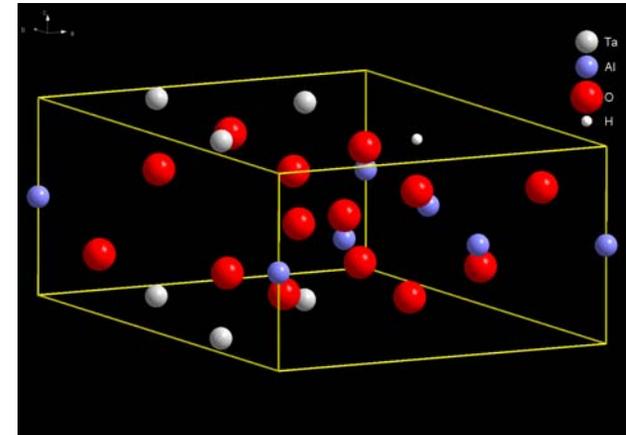
observed intensity modulation

An “Easy” Demonstration Sample



Simpsonite*

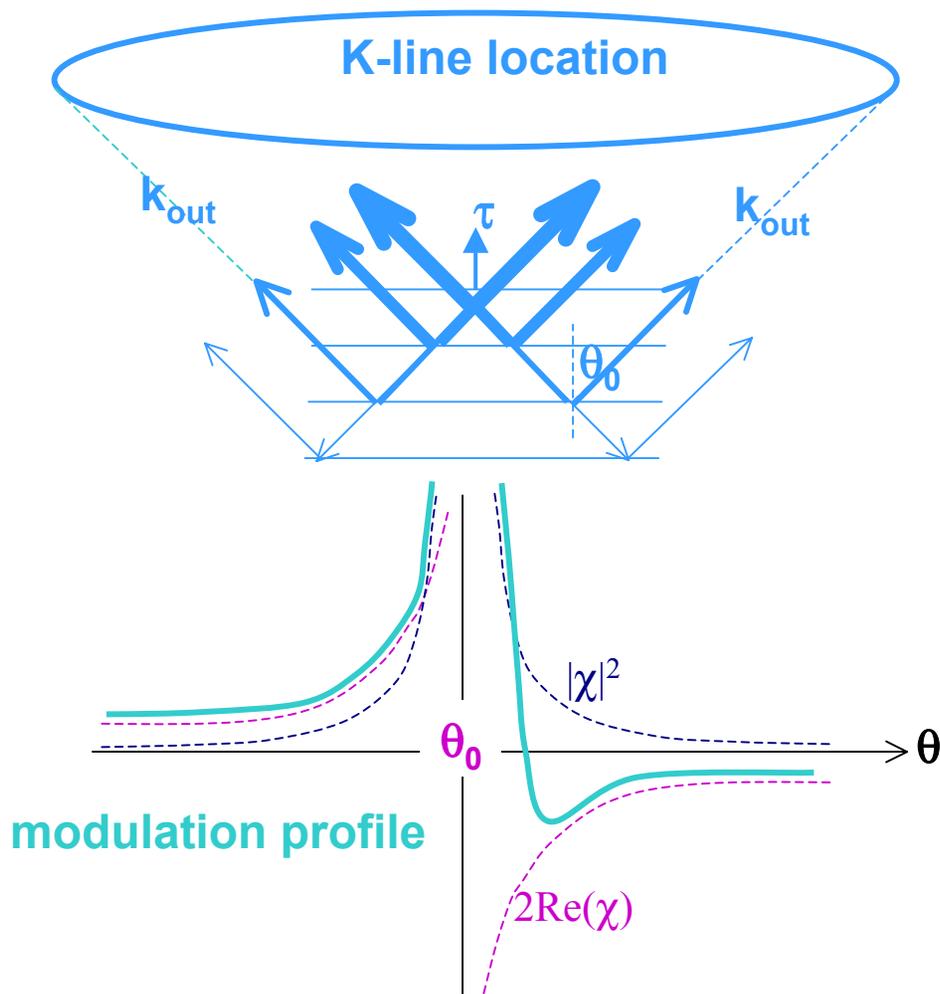
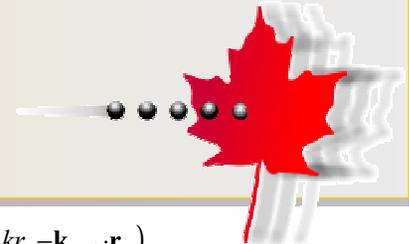
- No connection to Homer or Bart
- Mineral first discovered in Australia
- Simple structure
- 1 H atom



* Thanks to Ian Swainson (NRC) and Scott Ecrit (Canadian Museum of Nature)

Illustrations courtesy of Rob Hammond

s-Wave diffraction from planes: "K-Lines"



$$\chi(\mathbf{k}_{out}) = \int_{r_o} \frac{b(\mathbf{r}_o) e^{i(kr_o - \mathbf{k}_{out} \cdot \mathbf{r}_o)}}{r_o} d\mathbf{r}_o$$

$$B(\mathbf{q}) = \frac{1}{\sqrt{(2\pi)^3}} \int b(\mathbf{r}) e^{-i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}$$

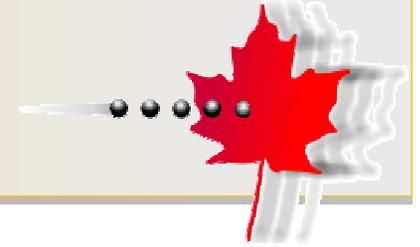
$$\chi(\mathbf{k}_{out}) = \frac{4\pi}{\sqrt{(2\pi)^3}} \int \frac{B(\mathbf{q})}{|\mathbf{q}|^2 - 2\mathbf{k}_{out} \cdot \mathbf{q}} d\mathbf{q}$$

$$B(\mathbf{q}) = \sum_{h,k,l} F_{hkl} \delta(\mathbf{q} - h\boldsymbol{\tau}_a - k\boldsymbol{\tau}_b - l\boldsymbol{\tau}_c)$$

$$\chi(\mathbf{k}_{out}) = \frac{4\pi}{\sqrt{(2\pi)^3}} \sum_{h,k,l} \frac{F_{hkl}}{|\boldsymbol{\tau}_{hkl}|^2 - 2\mathbf{k}_{out} \cdot \boldsymbol{\tau}_{hkl}}$$

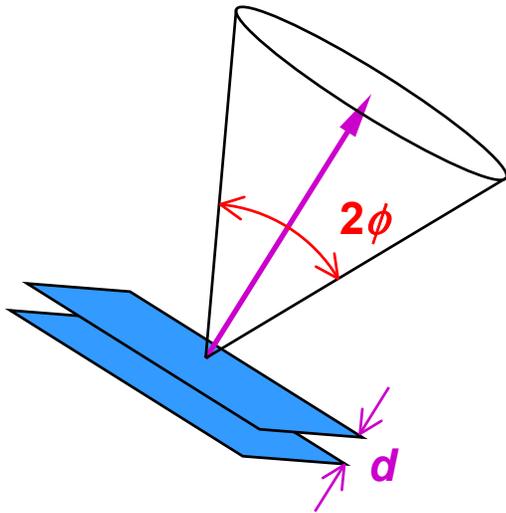
$$I_{1s} \cong \frac{a^* a}{R^2} \left[1 + 2\text{Re}(\chi(\mathbf{k}_{out})) + |\chi(\mathbf{k}_{out})|^2 \right]$$

What We Expect to See



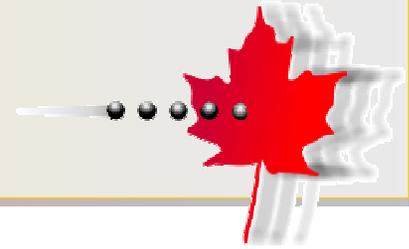
Interference leads to cones of scattering with:

- cone axis parallel to the plane normal
- characteristic angle relates wavelength and atomic spacing

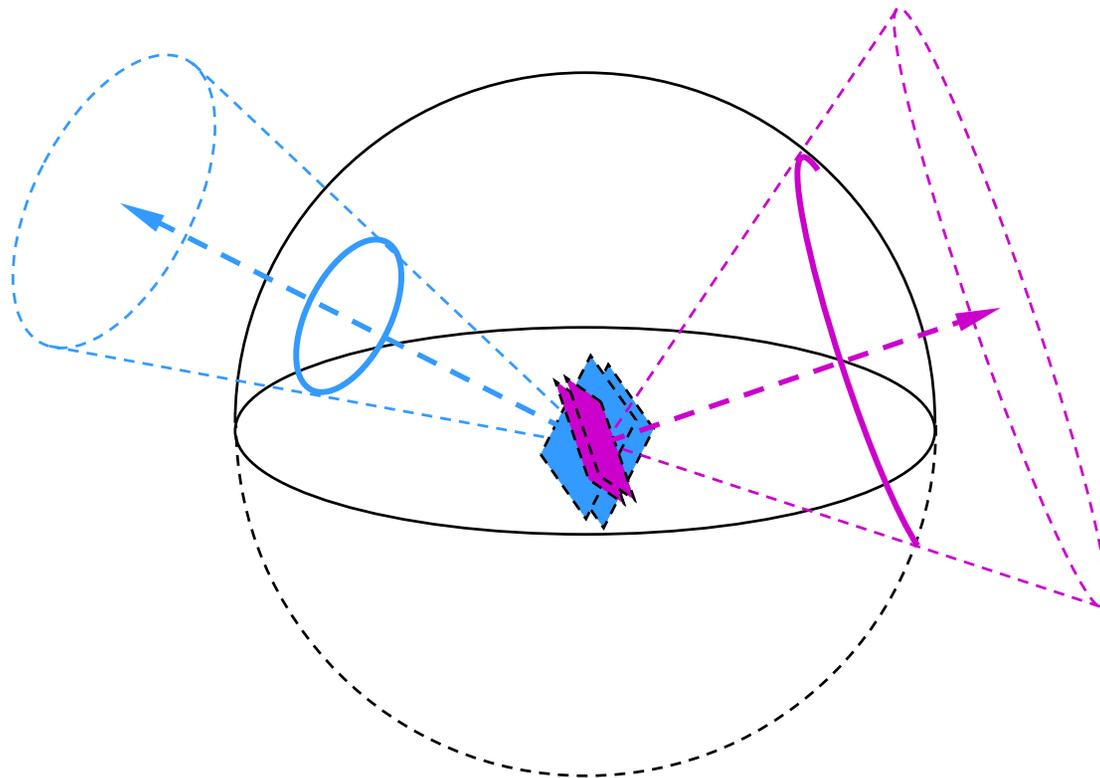


$$\lambda = 2d \cos \phi$$

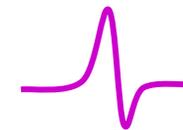
Expect to See “K-Lines”



Looking over a sphere of scattering, these cones intersect the sphere forming rings, known as Kossel Lines, Kikuchi Lines or collectively as “K-Lines”

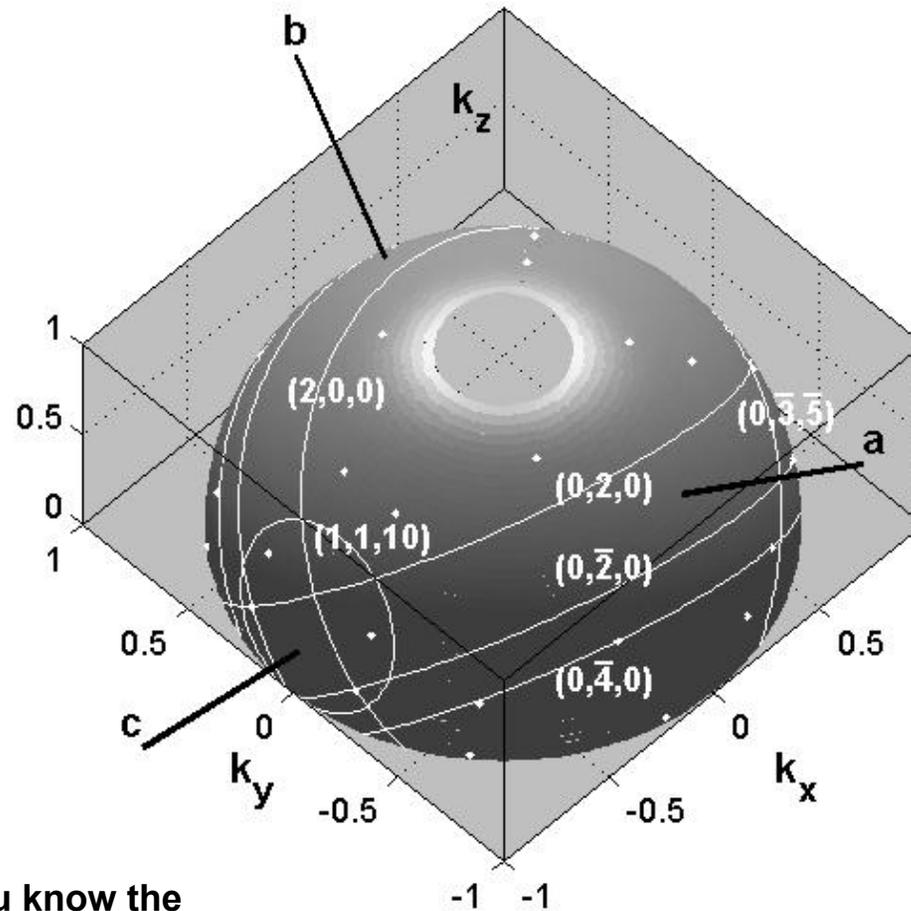
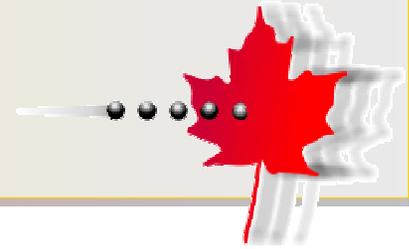


Intensity variations across the K-lines



Due to symmetry, a hemisphere is sufficient

Expect to See “K-Lines”

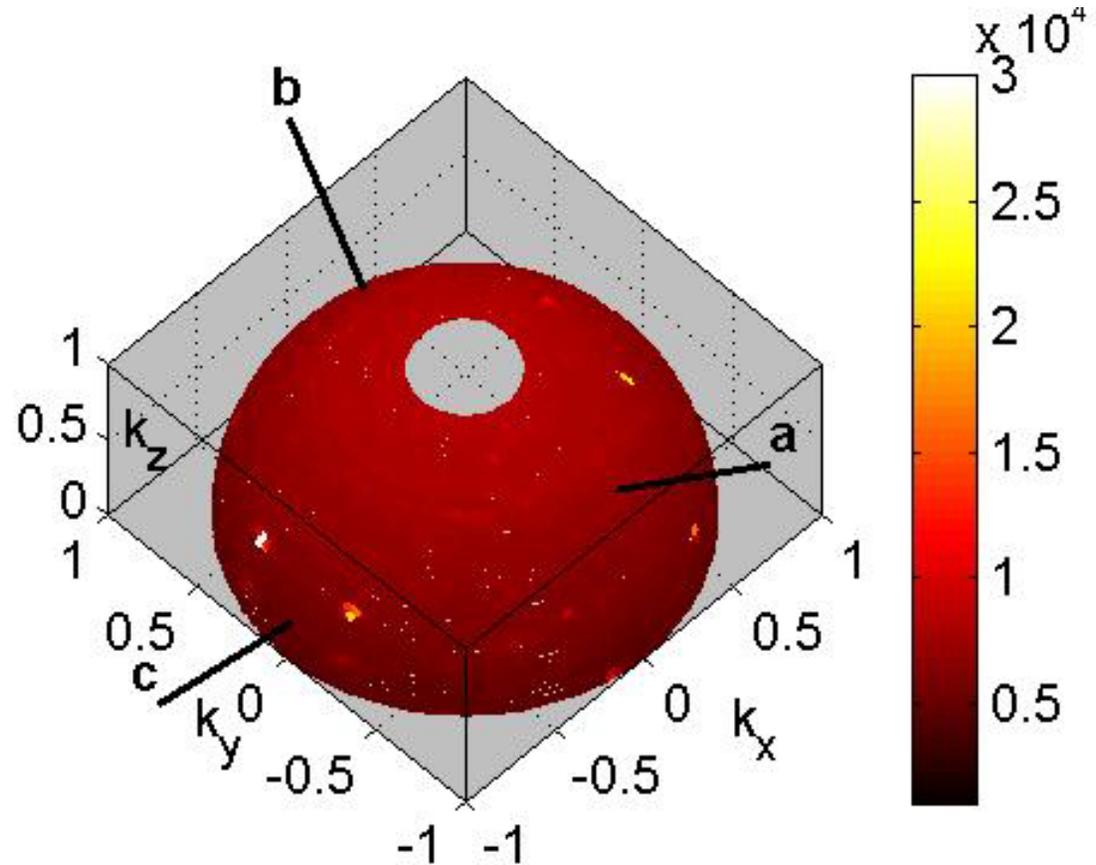


Can easily calculate once you know the crystal orientation and radiation wavelength.

K-lines: A Needle in a Hay Stack



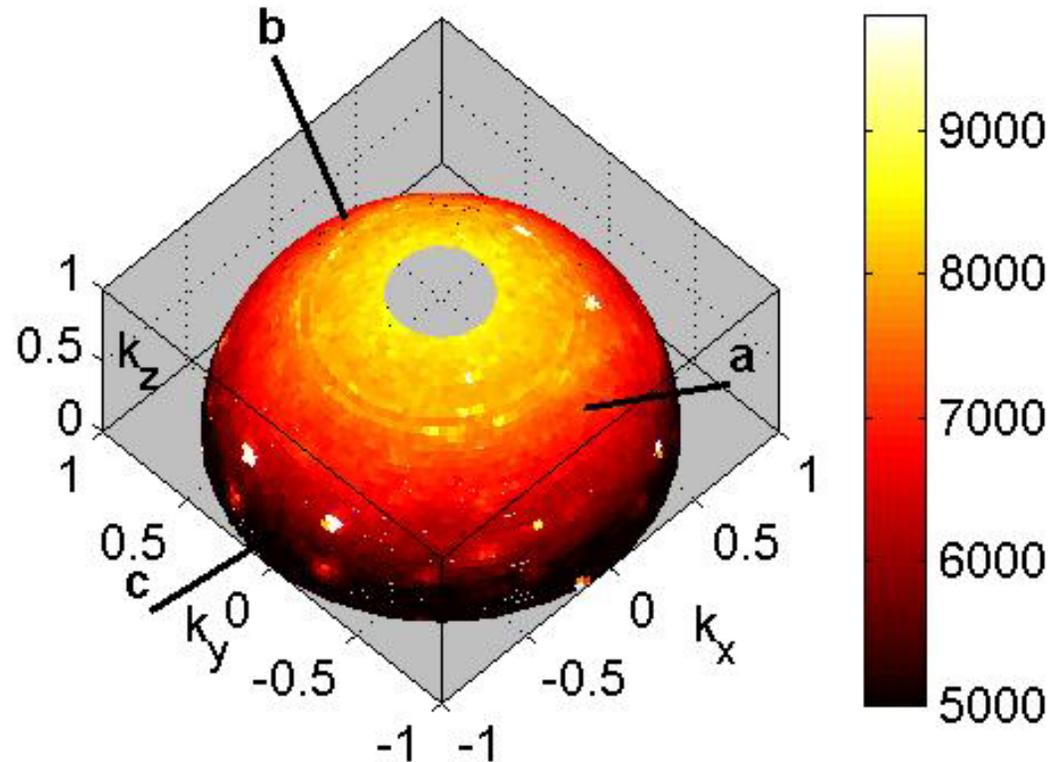
- Bragg peaks very intense
- Hologram data is in the featureless “background”



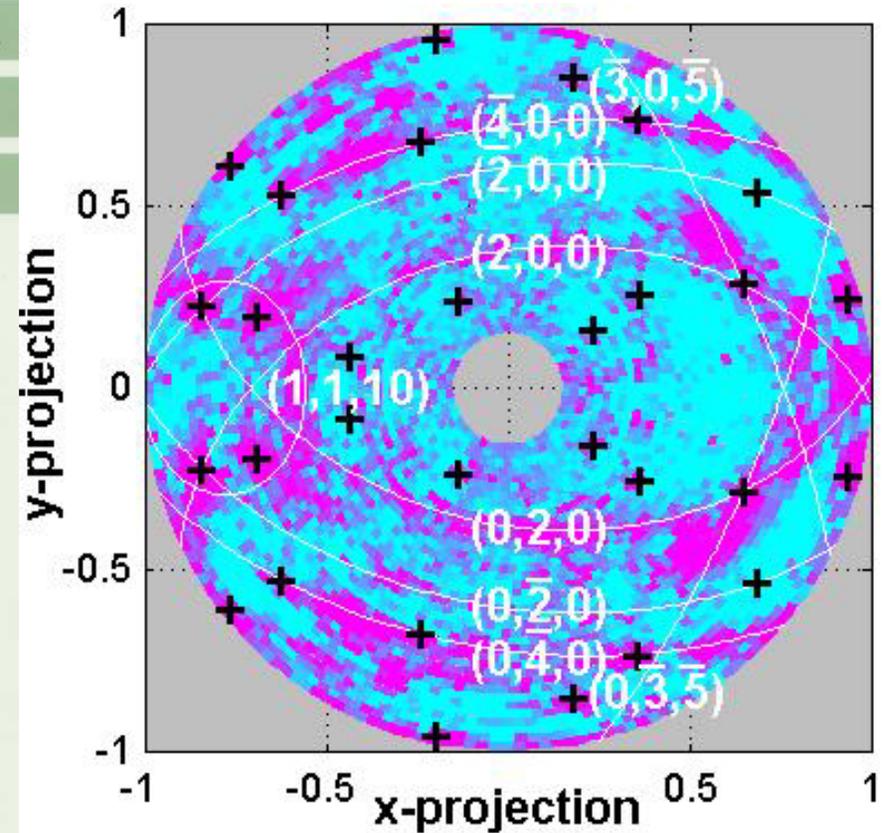
Looking for the Needle Reduce Dynamic Range



- The “background” is not featureless, but still dominated by slowly varying intensity
 - Attenuation
 - Debye-Waller factor
- Remove by fitting to slowly varying functions in ϕ and θ



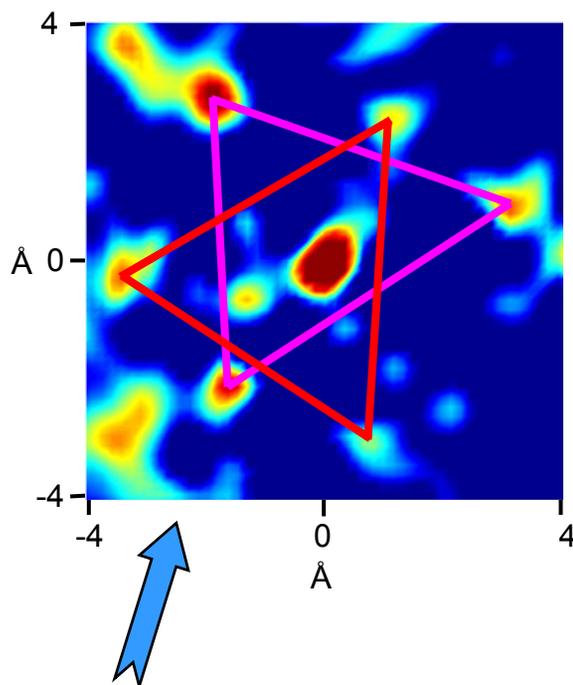
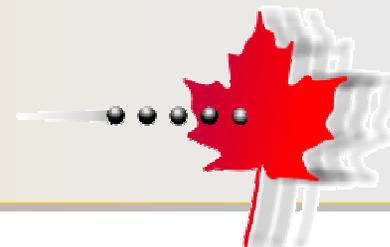
Found the Needle K-Lines Appear



Observation of Kossel and Kikuchi Lines in Thermal Neutron Incoherent Scattering
 B. Sur, R. B. Rogge, R. P. Hammond, V. N. P. Anghel, and J. Katsaras
 PRL, **88** (6), 065505 (2002)

NRC - CNRC

The Reconstruction

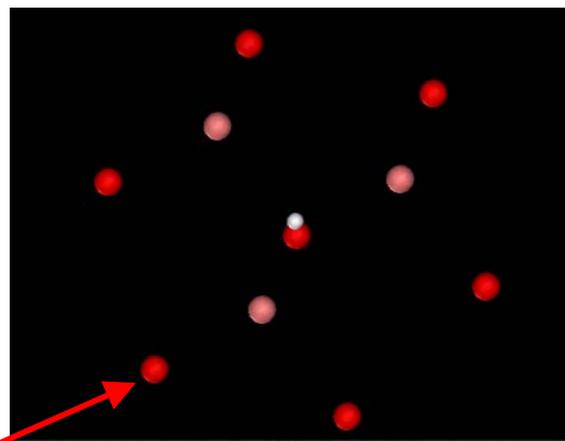


Slice of 3D data set taken at the known distance *below* the H atom to the plane of **red** O atoms

No model system required

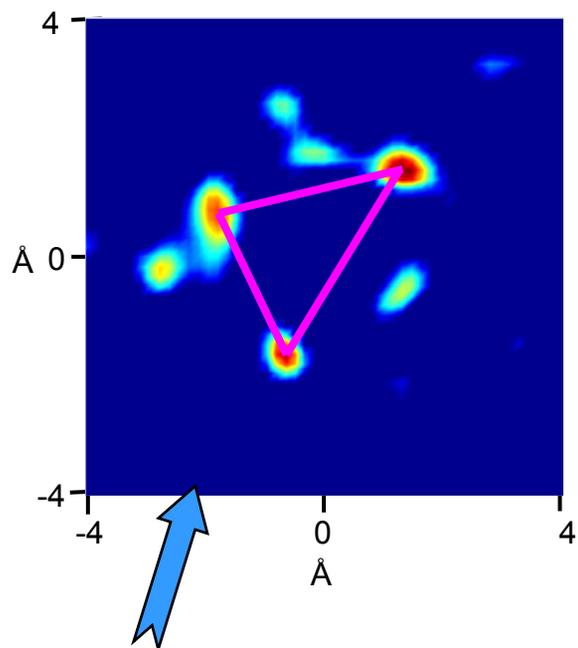
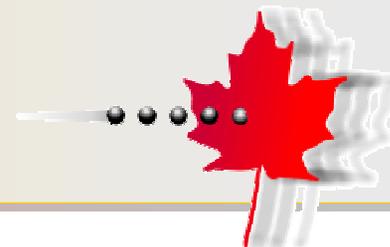
Crystal structure falls out of reconstruction

Obtain a full 3D view (H atom view)

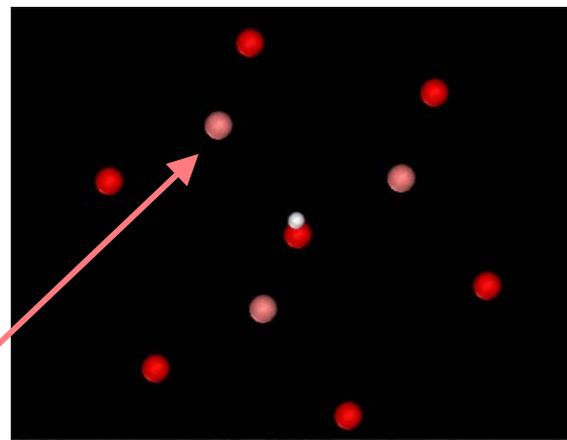


Atomic Structure Holography Using Thermal Neutrons
B. Sur, R. B. Rogge, R. P. Hammond, V. N. P. Anghel and J. Katsaras
Nature **414**, 525 (2001)

The Reconstruction



Slice from 3D data set taken at the known distance *above* the H atom to the plane of *light red* O atoms



Applications



- **Hologram of atoms in general... ***
 - **Solves the phase problem**
 - **Permits study of systems that cannot be crystallized (e.g. biological macromolecules, catalyst particles, nanostructures);**
 - **Has been used to:**
 - **locate dopant atoms in semiconductors and alloys;**
 - **probe ordering in alloys, quasicrystals;**
 - **look at interfaces and films**

* J. Spence, Nature (News and Views), 410, 1037-1040 (2001)

Applications Neutron Holography



- Neutrons are coherently scattered equally well by light and heavy atoms
- Neutrons can distinguish similar-Z atoms and can distinguish some isotopes
- The internal source H atom is abundant in biological materials and polymers
- H and D scatter very differently
- Neutrons are gentle, causing little or no damage to delicate systems

Take advantage of the complementarity of neutrons to tackle those systems that cannot be easily studied with electromagnetic radiation