

High Energy Physics with Ultra Cold Neutrons

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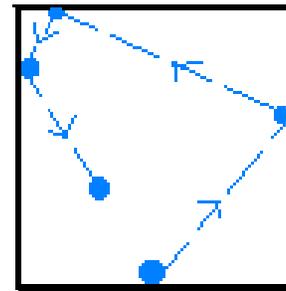
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- ❑ Advantageous properties of Ultra Cold Neutrons (UCN).
- ❑ A few possible experiments with UCN and current/future UCN sources.
- ❑ Focus: Measurements of angular correlations.
- ❑ Recent progress on understanding and measuring UCN depolarization.

UCN Are very Slow

$$E \leq 10^{-7} eV, \quad v \leq 8m/s, \quad t \leq 1mK$$

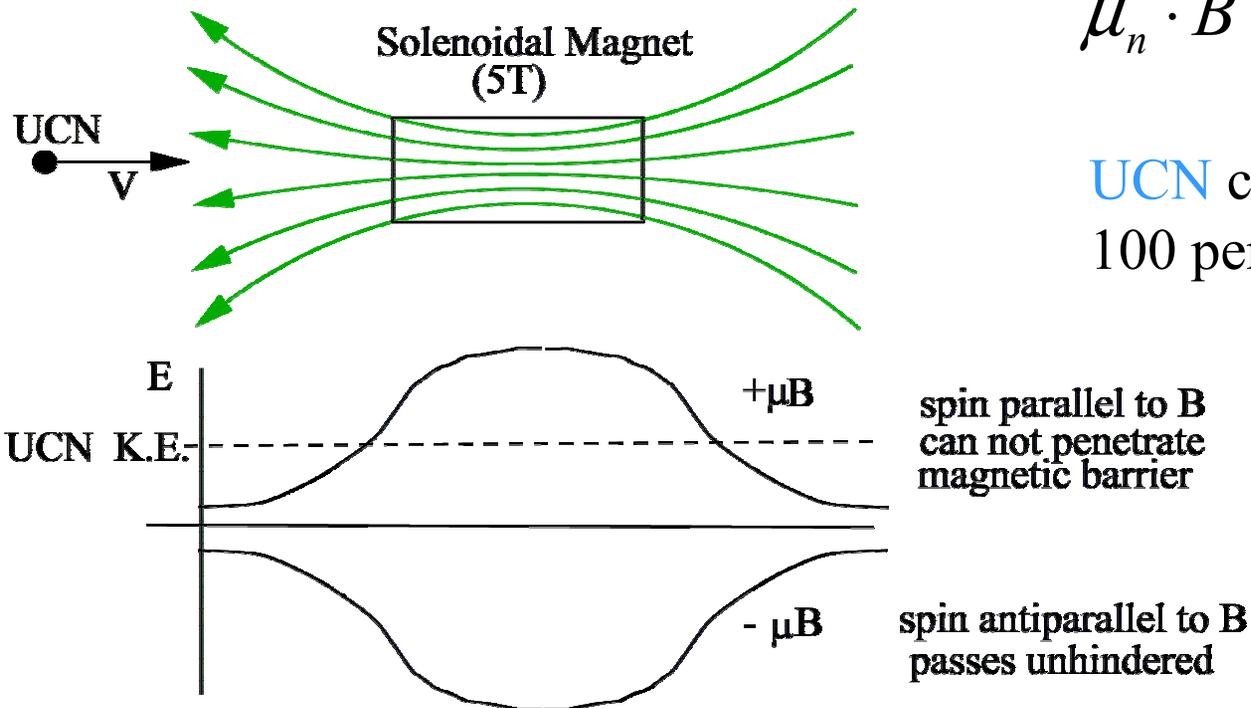
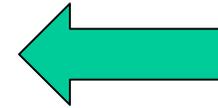


- Their wavelength is long ($\sim 500\text{\AA}$), thus, they totally internally reflect from many surfaces. (Ex: ^{58}Ni , Diamond,...).
 - They can be guided long distances or stored in bottles for long times.
 - Neutron experiments can be undertaken in a well shielded, low radiation background environment at reactors or (especially) spallation sources.
 - The UCN activation of an experimental apparatus is negligible.

UCN are Easy to Polarize



UCN kinetic energy is less than their potential energy in a strong magnetic field.



$$\mu_n \cdot B \approx 3 \times 10^{-7} \text{ eV}$$

UCN can be essentially 100 percent polarized.

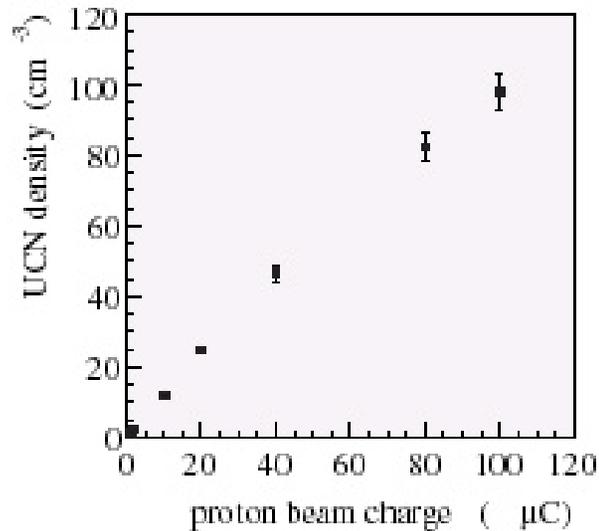
(note: neutron magnetic moment is negative)

Experiments with UCN

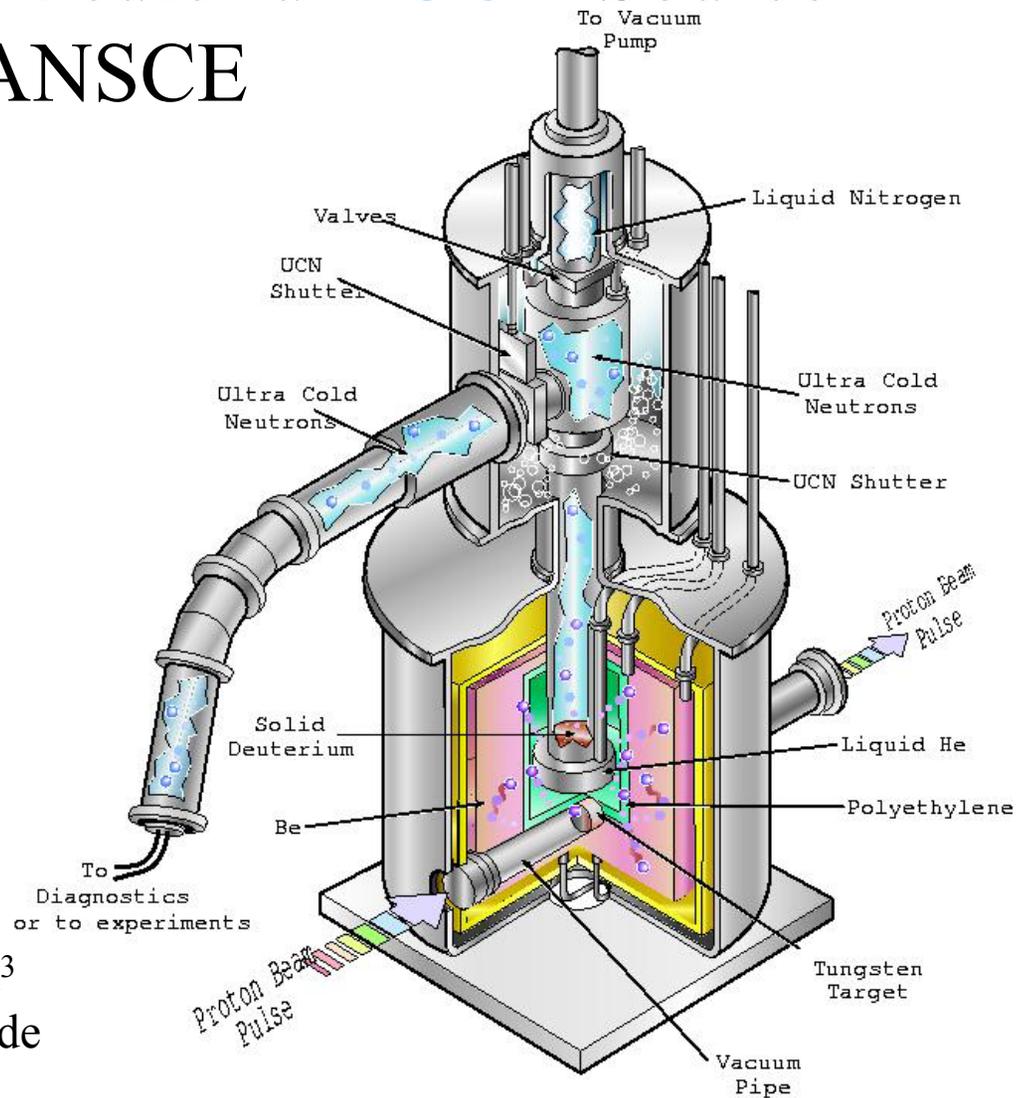
- High polarization, low radiation background and long storage time facilitate many high precision measurements of neutron properties.
- For example:
 - The neutron lifetime.
 - The neutron electric dipole moment.
 - Neutron-antineutron oscillation.
 - The beta energy spectrum.
 - The recoil proton energy spectrum.
 - The angular correlations between neutron decay products (electron, proton and neutron spin).
- These quantities are interesting in themselves. Perhaps more intriguing, they are also probes of “new” physics.
- Recent advances in UCN production make efforts feasible.

A Superthermal Solid Deuterium UCN Source at LANSCE

World record densities achieved

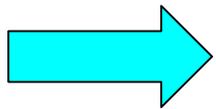


Compare to previous record of 41 UCN/cm³ (at ILL). Note: over two orders of magnitude improvement may ultimately be possible.



Current and Future World Wide UCN Sources

Institution	Source Type	Density (UCN/cc)
ILL	Reactor	40
LANL	Spallation	400-800
FRMII	Reactor	10^4
PSI	Spallation	2.5×10^3
NCSU	Reactor	1.5×10^3



Operational for
experiment,
Fall 2003.

Angular Correlations of Neutron Decay: The Current Situation

$$WdE_e d\Omega_e d\Omega_\nu \propto p_e E_e (E_0 - E_e)^2 \left(1 + a \frac{\vec{p}_e \cdot \hat{p}_\nu}{E_e} + b \frac{m_e}{E_e} + A \vec{\sigma} \cdot \frac{\vec{p}_e}{E_e} + B \vec{\sigma} \cdot \hat{p}_\nu + D \vec{\sigma} \cdot \frac{\vec{p}_e \times \hat{p}_\nu}{E_e} \right)$$

To first order in the SM:

$$a = \frac{1 - \lambda^2}{1 + 3\lambda^2}$$

$$A = -2 \frac{\lambda(\lambda + 1)}{1 + 3\lambda^2}$$

$$B = 2 \frac{\lambda(\lambda - 1)}{1 + 3\lambda^2}$$

$$\lambda = g_a / g_v$$

	PDG 2000
a	-0.102±0.005
A	-0.1162±0.0013
B	0.983±0.004
D	(-0.5 ±1.4)×10 ⁻³
λ	-1.2670±.0035

Note: A **UCN** measurement of **D** may have less systematic error than a **CN** measurement.

Physics motivation for measuring λ with UCN

- There is a $2\text{-}\sigma$ discrepancy between V_{ud} determined from $0^+ \rightarrow 0^+$ nuclear beta decay (most precise measurement to date) and CKM unitarity.
- Neutron decays are amenable to a straightforward theoretical treatment because they are a single nucleon, (as opposed to $0^+ \rightarrow 0^+$ decays), and thus, the best way to learn more.
- Given λ , the neutron lifetime and the muon lifetime, one can determine V_{ud} .
- A is the most sensitive of the three correlations to λ :

$$\frac{d\lambda}{dA} = 2.6, \quad \frac{d\lambda}{da} = 3.3, \quad \frac{d\lambda}{dB} = 13.4$$

- Measuring B with the same experimental apparatus provides a polarization independent determination of λ . (I.A. Kuznetsov et al. NIM A 440 (2000) 539-542).

Additional Motivations for **a** and **B**

- **A**, **a** and **B** are sensitive to different recoil order corrections. (Example: **a** is most sensitive to weak magnetism and induced tensor form factors).
- A precise measurements of **A** and **a** at the 0.1% level permits a test of the conserved-vector current hypothesis (CVC) and a search for second class currents (SCC). (S. Gardner and C. Zhang PRL 86(2001) p.5666).
- A precise measurement constrains both the minimally super symmetric model (MSSM) and the R-violating extension to MSSM. (Ramsey-Musolf PRL 88(2002) 071804, PRD 62(2000) 056009).
- A precise measurement of **B** constrains $\left(\frac{M_L}{M_R}\right)^2$ in manifest left-right symmetric models. (Kuznetsov et al. PRL 75(1995) p. 794).

The UCN A Collaboration

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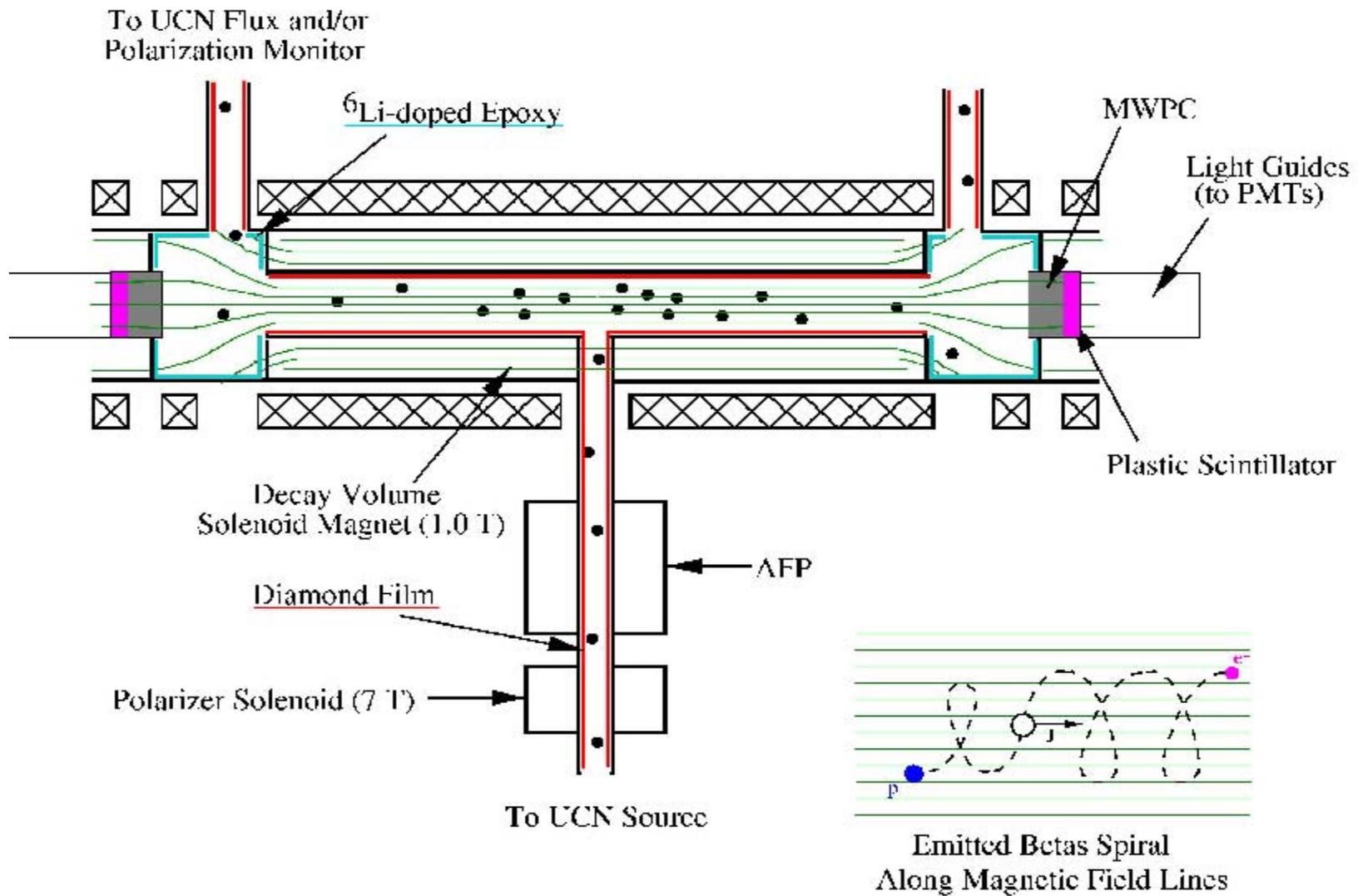
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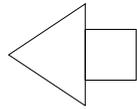
Initial goal: $\Delta A/A = 0.2\%$

A Spectrometer for Angular Correlation Studies



How to Measure A : The Electron Asymmetry

1

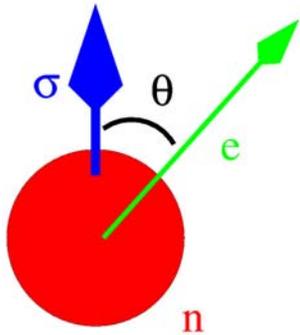


Neutron Spin

2

$$dW = [1 + \beta P A(E) \cos \theta] d\Gamma(E)$$

$$A_{\text{exp}}(E) \propto \frac{N^{\uparrow}(E) - N^{\downarrow}(E)}{N^{\uparrow}(E) + N^{\downarrow}(E)} \approx \frac{1}{2} A(E) \beta P$$

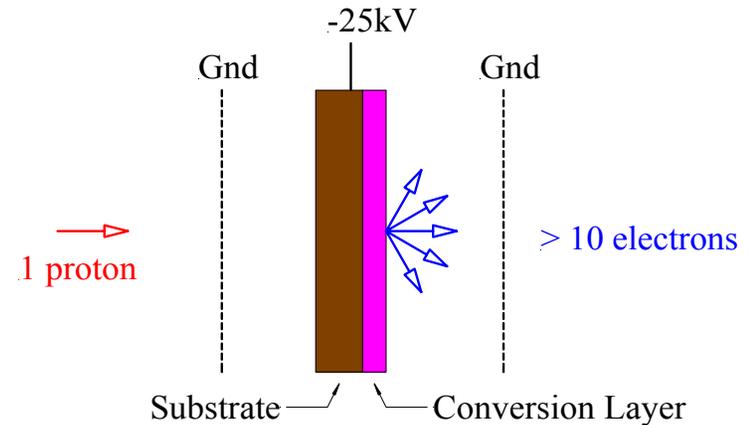


- Measure electron energy, electron direction, neutron polarization.
- Controlling systematic errors is critical:
 - Background subtraction.
 - Polarization (0.1%)
 - Detector Response (backscatter, linearity etc.) (0.1%)

We are undertaking a thorough study to understand and control systematic errors associated with beta detectors (MWPC, Scin., Si.).

Electron Transparent Proton Detectors

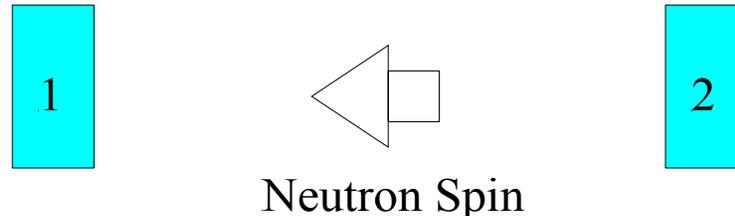
- The Strategy: Convert protons into electrons in an electron transparent foil which has no impact on A . Use the planned electron detectors to see these electrons.
- Advantages:
 - Should be easy to install, easy to operate, robust proton identification.
 - Minimal impact on systematic errors of A measurement (electron backscatter and energy loss).
 - Measure B concurrently with A .



- PERKEO colleagues tried C-foils ($25\mu\text{g}/\text{cm}^2$) with MgO.
- We are investigating fluorinated polyimide ($0.1\ \mu\text{m}$) with CsI or diamond (200\AA).
- Should be stronger, larger, with a higher electron yield.

How to Measure **B**: The Proton Asymmetry

- The proton asymmetry is the best method in our geometry.
- Measure emission direction of protons only (electron coincidence not required).



$$\alpha_p = \frac{p_+ - p_-}{p_+ + p_-} = PC_p(A + B), \quad C_p = 0.27484$$

Polarization uncertainty limited measurement with 10^7 decays (36hrs).

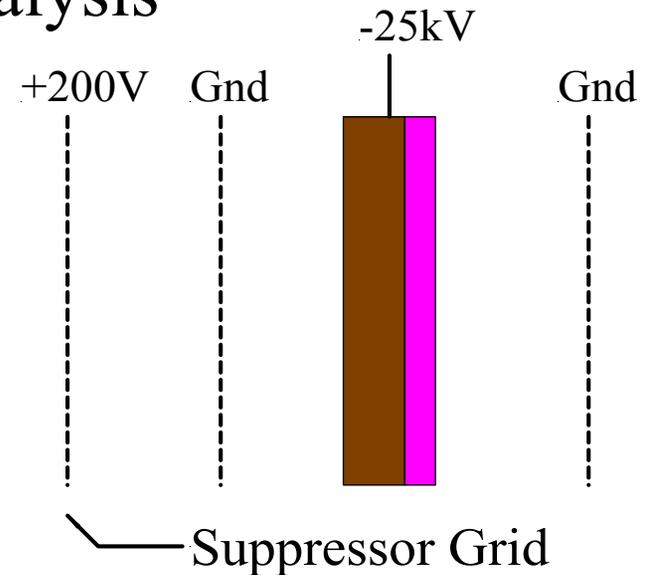
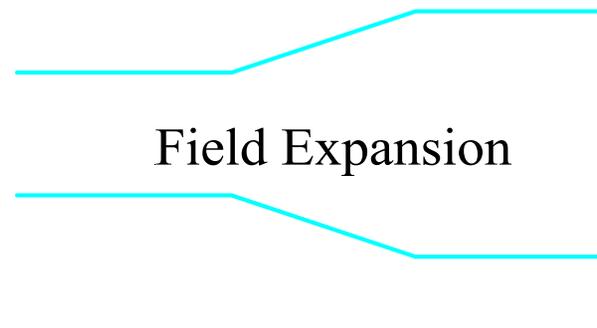
$$\Delta B/B = 0.2\%$$

Gluck et al. *Nuclear Physics A* 593 (1995) pg. 125

How to Measure a : The 1-d Integration Method

A Preliminary Analysis

- 1) Fill trap and count protons higher than E_{cut} .
- 2) Empty trap and count protons lower than E_{cut} .

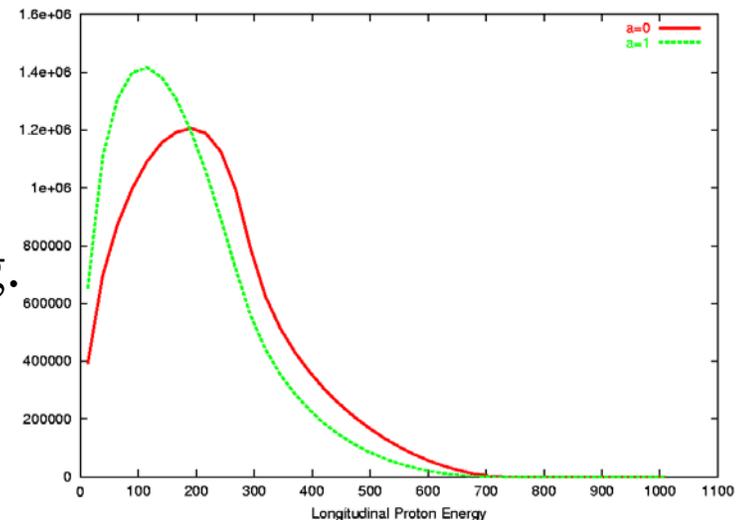


$\Delta a/a = 0.5\%$

5×10^7 decays
(8 days)

Very preliminary investigation.
Systematic errors are very challenging.
Field uniformity, unpolarization, etc.

J. Byrne et al. NIMA 349(1994) pg. 454



Polarization Progress

Polarization Level

- Goal for the first experiments: less than 1×10^{-3} depolarization error (we should do significantly better).

Sources of Depolarization

- Material depolarization – Benchmarked at the ILL (fall 2001, spring 2002).
 - less than 2×10^{-6} per bounce
 - 8×10^{-5} per sec
- Majorana transitions – Monte Carlo treatment complete.
 - less than 2×10^{-4} per pass (holding fields $\leq 200\text{G}$)
 - in field reversal region
- Wall collisions in gradient fields - Monte Carlo exists.
 - less than 1×10^{-4} per pass in field reversal region and AFP region
- AFP performance – Monte Carlo exists, partially benchmarked at the ILL.
 - less than 1×10^{-4} per pass